

Poisson Bracket:

These are quantities which are invariant under canonical transformations and are the classical analogues of commutation relations between operators in quantum mechanics.

The Poisson bracket of two fns. u, v w.r.t the canonical variables (q, p) is defined as:

$$[u, v]_{q, p} = \sum_{i=1}^n \left(\frac{\partial u}{\partial q_i} \frac{\partial v}{\partial p_i} - \frac{\partial u}{\partial p_i} \frac{\partial v}{\partial q_i} \right)$$

Usually the Einstein convention is used and $[u, v]_{q, p} = [u, v] = \frac{\partial u}{\partial q_i} \frac{\partial v}{\partial p_i} - \frac{\partial u}{\partial p_i} \frac{\partial v}{\partial q_i}$

(3) If not specifically mentioned, it is understood that the P.B are evaluated w.r.t (q, p)

Some properties:

(1) $[u, v] = -[v, u]$

(2) $[u, u] = [v, v] = 0$

(3) $[cu, v] = [u, cv] = c[u, v]$

for c - ~~not~~ constant $[cu, v] = \frac{\partial (cu)}{\partial q} \frac{\partial v}{\partial p} - \frac{\partial (cu)}{\partial p} \frac{\partial v}{\partial q}$

$$\begin{aligned}
 (3) \quad [u+v, w] &= \frac{\partial}{\partial q_i} (u+v) \frac{\partial w}{\partial p_i} - \frac{\partial}{\partial p_i} (u+v) \frac{\partial w}{\partial q_i} \\
 &= \left(\frac{\partial u}{\partial q_i} + \frac{\partial v}{\partial q_i} \right) \frac{\partial w}{\partial p_i} - \left(\frac{\partial u}{\partial p_i} + \frac{\partial v}{\partial p_i} \right) \frac{\partial w}{\partial q_i} \\
 &= \frac{\partial u}{\partial q_i} \frac{\partial w}{\partial p_i} + \frac{\partial v}{\partial q_i} \frac{\partial w}{\partial p_i} - \frac{\partial u}{\partial p_i} \frac{\partial w}{\partial q_i} - \frac{\partial v}{\partial p_i} \frac{\partial w}{\partial q_i}
 \end{aligned}$$

Re-arranging terms we get,

$$[u+v, w] = [u, w] + [v, w] \text{ — distributive property.}$$

$$(4) \quad [u, vw] = [u, v] w + v [u, w]$$

$$(5) \quad \frac{\partial}{\partial t} [u, v] = \left[\frac{\partial u}{\partial t}, v \right] + \left[u, \frac{\partial v}{\partial t} \right]$$

(6) Jacobi's identity:

$$[u, [v, w]] + [v, [w, u]] + [w, [u, v]] = 0$$

$$[u, [v, w]] + [w, [u, v]] + [v, [w, u]] = 0$$

$$(7) \quad [u, F(w_1, w_2)] = \frac{\partial u}{\partial q_i} \frac{\partial F}{\partial w_1} [u, w_1] + \frac{\partial F}{\partial w_2} [u, w_2]$$

Proof.

$$\begin{aligned}
 [u, F(w_1, w_2)] &= \frac{\partial u}{\partial q} \frac{\partial}{\partial p} (F(w_1, w_2)) - \frac{\partial u}{\partial p} \frac{\partial}{\partial q} (F(w_1, w_2)) \\
 &= \frac{\partial u}{\partial q} \left(\frac{\partial F}{\partial w_1} \frac{\partial w_1}{\partial p} + \frac{\partial F}{\partial w_2} \frac{\partial w_2}{\partial p} \right) - \frac{\partial u}{\partial p} \left(\frac{\partial F}{\partial w_1} \frac{\partial w_1}{\partial q} + \frac{\partial F}{\partial w_2} \frac{\partial w_2}{\partial q} \right) \\
 &= \frac{\partial F}{\partial w_1} \left(\frac{\partial u}{\partial q} \frac{\partial w_1}{\partial p} - \frac{\partial u}{\partial p} \frac{\partial w_1}{\partial q} \right) + \frac{\partial F}{\partial w_2} \left(\frac{\partial u}{\partial q} \frac{\partial w_2}{\partial p} - \frac{\partial u}{\partial p} \frac{\partial w_2}{\partial q} \right) \\
 &= \frac{\partial F}{\partial w_1} [u, w_1] + \frac{\partial F}{\partial w_2} [u, w_2]
 \end{aligned}$$

In general, for $F = (w_1, w_2, \dots, w_n)$

$$\begin{aligned}
 [u, F(w_1, w_2, \dots, w_n)] &= \frac{\partial F}{\partial w_1} [u, w_1] + \frac{\partial F}{\partial w_2} [u, w_2] \\
 &+ \frac{\partial F}{\partial w_3} [u, w_3] + \dots + \frac{\partial F}{\partial w_n} [u, w_n]
 \end{aligned}$$

Elementary P.B.

P.B.'s constructed out of the canonical variables themselves (co-ord & momenta) are called elementary P.B.'s.

$$[q_i, p_j] = \frac{\partial q_i}{\partial q_i} \frac{\partial p_j}{\partial p_i} - \frac{\partial q_i}{\partial p_i} \frac{\partial p_j}{\partial q_i}$$

$$\frac{\partial q_i}{\partial q_i} = 1 \quad \frac{\partial p_j}{\partial p_i} = 1 \text{ only if } i=j$$

Properties:

$$(1) [q_i, q_j] = [p_i, p_j] = 0$$

$$(2) [q_i, p_j] = -[p_j, q_i] = \delta_{ij}$$

$$(3) [u, q_i] = -\frac{\partial u}{\partial p_i} \quad \& \quad [u, p_i] = \frac{\partial u}{\partial q_i}$$

$$(4) [u, \vec{r}] = \sum_i \left(\frac{\partial u}{\partial q_i} \frac{\partial \vec{r}}{\partial p_i} - \frac{\partial u}{\partial p_i} \frac{\partial \vec{r}}{\partial q_i} \right) \quad \text{--- (1)}$$

$$\text{Let } \vec{r} = x\hat{i} + y\hat{j} + z\hat{k}$$

Then the 2nd term on R.H.S of (1) is :

$$-\frac{\partial u}{\partial p_i} \sum_i \left(\hat{i} + \hat{j} + \hat{k} \right) = -\sum_i \left(\frac{\partial u}{\partial p_x} \hat{i} + \frac{\partial u}{\partial p_y} \hat{j} + \frac{\partial u}{\partial p_z} \hat{k} \right)$$

$$= -\vec{\nabla}_p u$$

$$(5) [u, \vec{p}] = \vec{\nabla}_r u$$

Hamilton's eqns in the PB formulation:

(1) For $u = u(q, p, t)$

$$\frac{du}{dt} = \sum_{i=1}^n \frac{\partial u}{\partial q_i} \dot{q}_i + \sum_{i=1}^n \frac{\partial u}{\partial p_i} \dot{p}_i + \frac{\partial u}{\partial t} \quad \text{--- (1)}$$

(2) But from the canonical eqns. we have.

$$\dot{q} = \frac{\partial H}{\partial p} \quad \& \quad -\dot{p} = \frac{\partial H}{\partial q}$$

$\therefore$ (1) becomes:

$$\frac{du}{dt} = \left(\frac{\partial u}{\partial q} \frac{\partial H}{\partial p} - \frac{\partial u}{\partial p} \frac{\partial H}{\partial q} \right) + \frac{\partial u}{\partial t}$$

$$\frac{du}{dt} = [u, H] + \frac{\partial u}{\partial t} \quad \text{--- (2)}$$

If, in (2), $u = q$, (2) becomes

$$\left. \begin{aligned} \dot{q} &= [q, H] \\ \dot{p} &= [p, H] \end{aligned} \right\} \begin{array}{l} \& \text{ when } u = p, \\ \text{--- Hamilton's equations} \\ \text{in p.B formulation} \end{array}$$

Also, if u is a constant of motion,

$$\frac{du}{dt} = 0 \quad \text{or} \quad [u, H] = - \frac{\partial u}{\partial t}$$

Jacobi-Poisson Theorem:

If u & v are any two constants of motion, then, their PB $[u, v]$ is also a constant of motion.

Since u & v are constants of motion, then,

$$\text{Given: } [u, H] = - \frac{\partial u}{\partial t} \quad \& \quad [H, u] = \frac{\partial u}{\partial t} \quad \&$$

$$[H, v] = \frac{\partial v}{\partial t}$$

To prove: $[H, [u, v]] = \frac{\partial}{\partial t} [u, v]$

Proof: By Jacobi's identity,

$$[w, [u, v]] + [v, [w, u]] + [u, [v, w]] = 0.$$

Putting $w = H$,

$$\Rightarrow [H, [u, v]] + [v, [H, u]] + [u, [v, H]] = 0$$

$$\begin{aligned} \text{or } [H, [u, v]] &= -[v, [H, u]] - [u, [v, H]] \\ &= [[H, u], v] + [u, [H, v]] \\ &= \left[\frac{\partial u}{\partial t}, v \right] + \left[u, \frac{\partial v}{\partial t} \right] = \frac{\partial}{\partial t} [u, v] \end{aligned}$$

ie if u & v are constants of motion,
then their P.B's are also constants
of motion.